

Artificial emissions and formation of ionized layers in the ionosphere due to HF-heating

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Motivations and Objectives

Models of electron acceleration due to HF-heating were introduced since 1970-es.

However recent optical and ionosonde observations emphasized a need of a comprehensive model of electron acceleration by the Langmuir turbulence due to the ionospheric HF-heating.

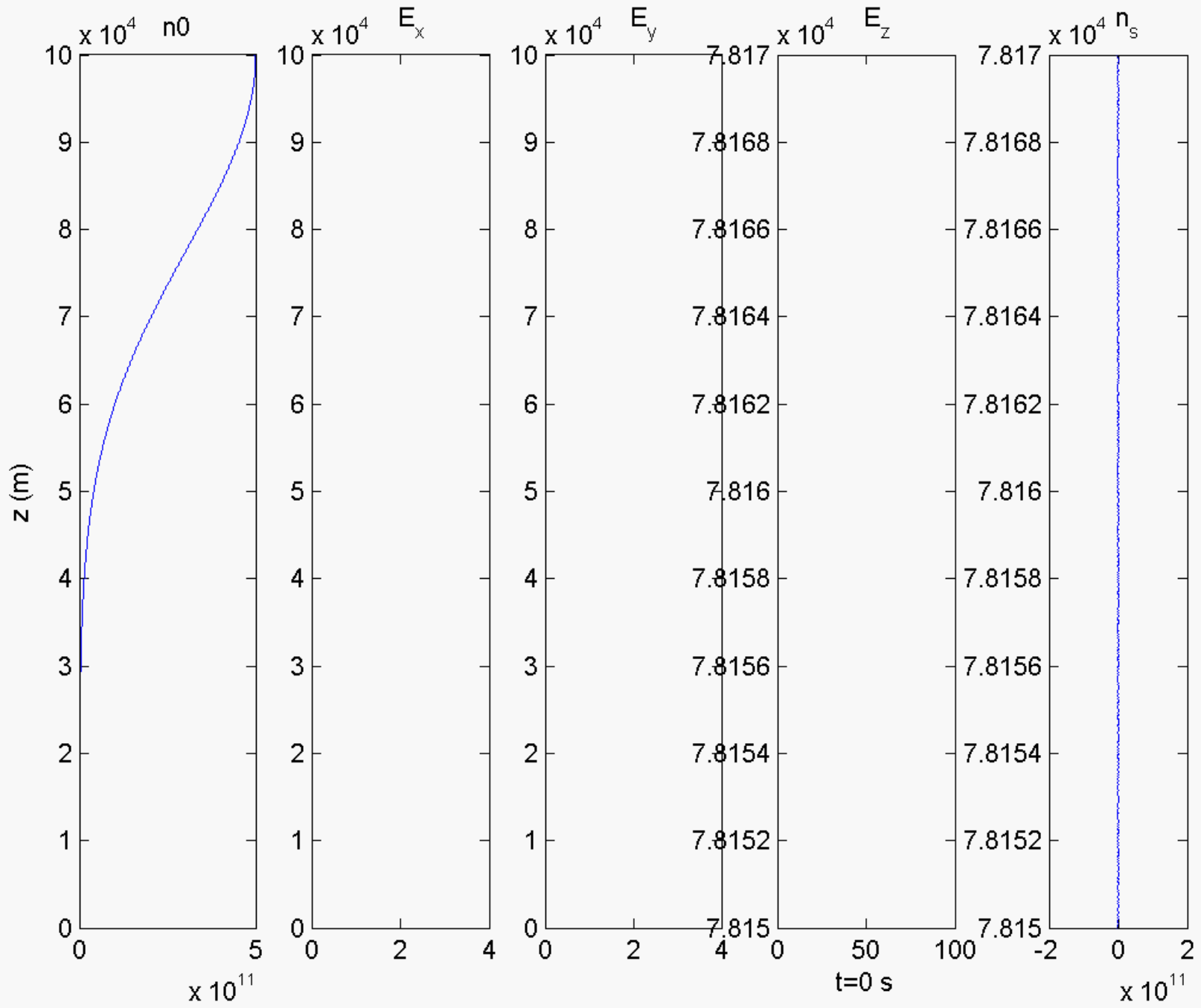
The objective of this project is to develop such a model and check it against the existing observations.

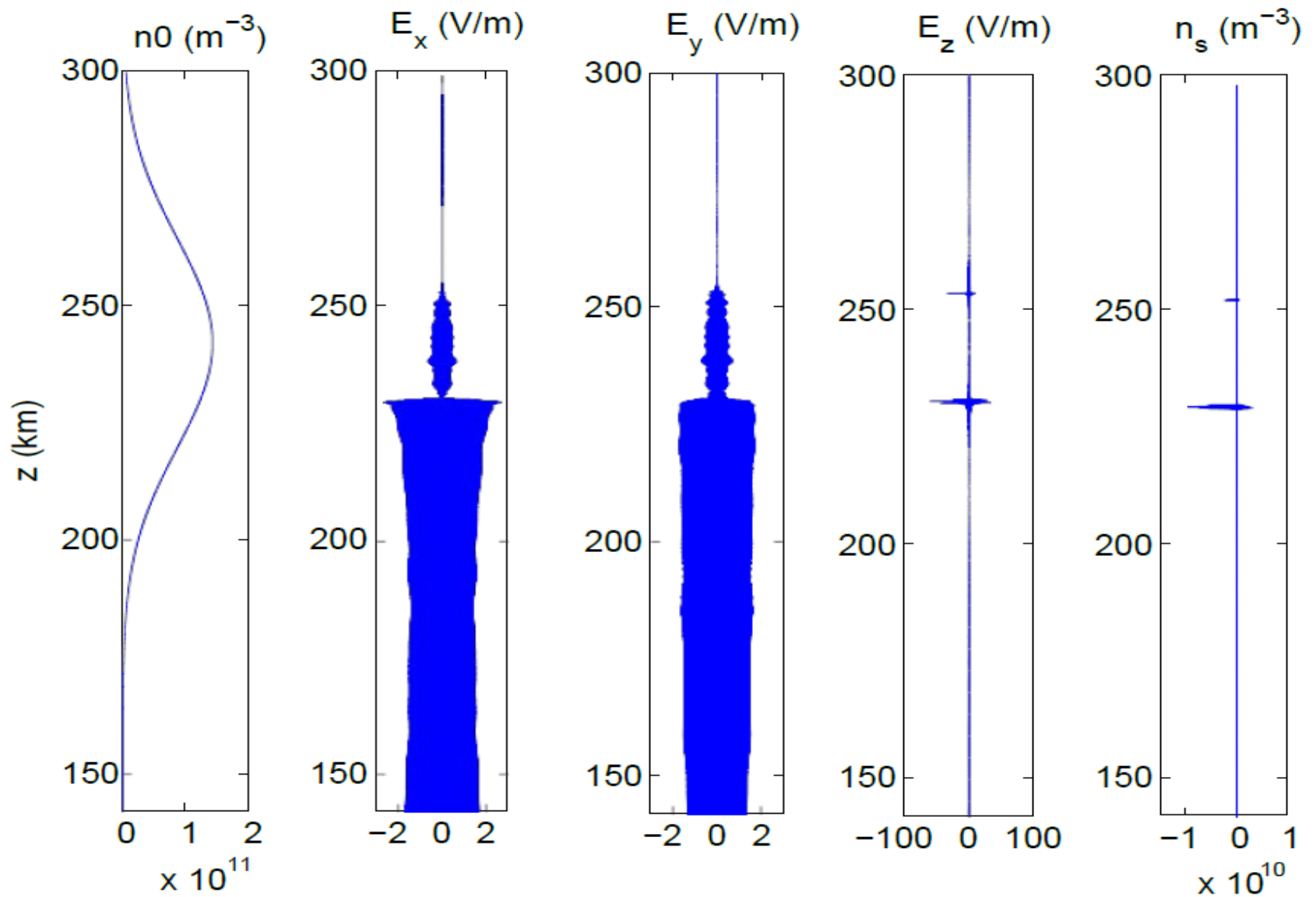
Model Approach

- Describe topside electrostatic turbulence due to HF-heating
- Obtain energy spectrum of the electrons accelerated by such turbulence:
 - Solving the F-P equation
 - Using PIC method
- Consider effects on the energy spectrum due to energy losses by the electrons propagating in the ionosphere
- Study effects caused by the fast electrons:
 - Optical emissions
 - Formation of ionized layers

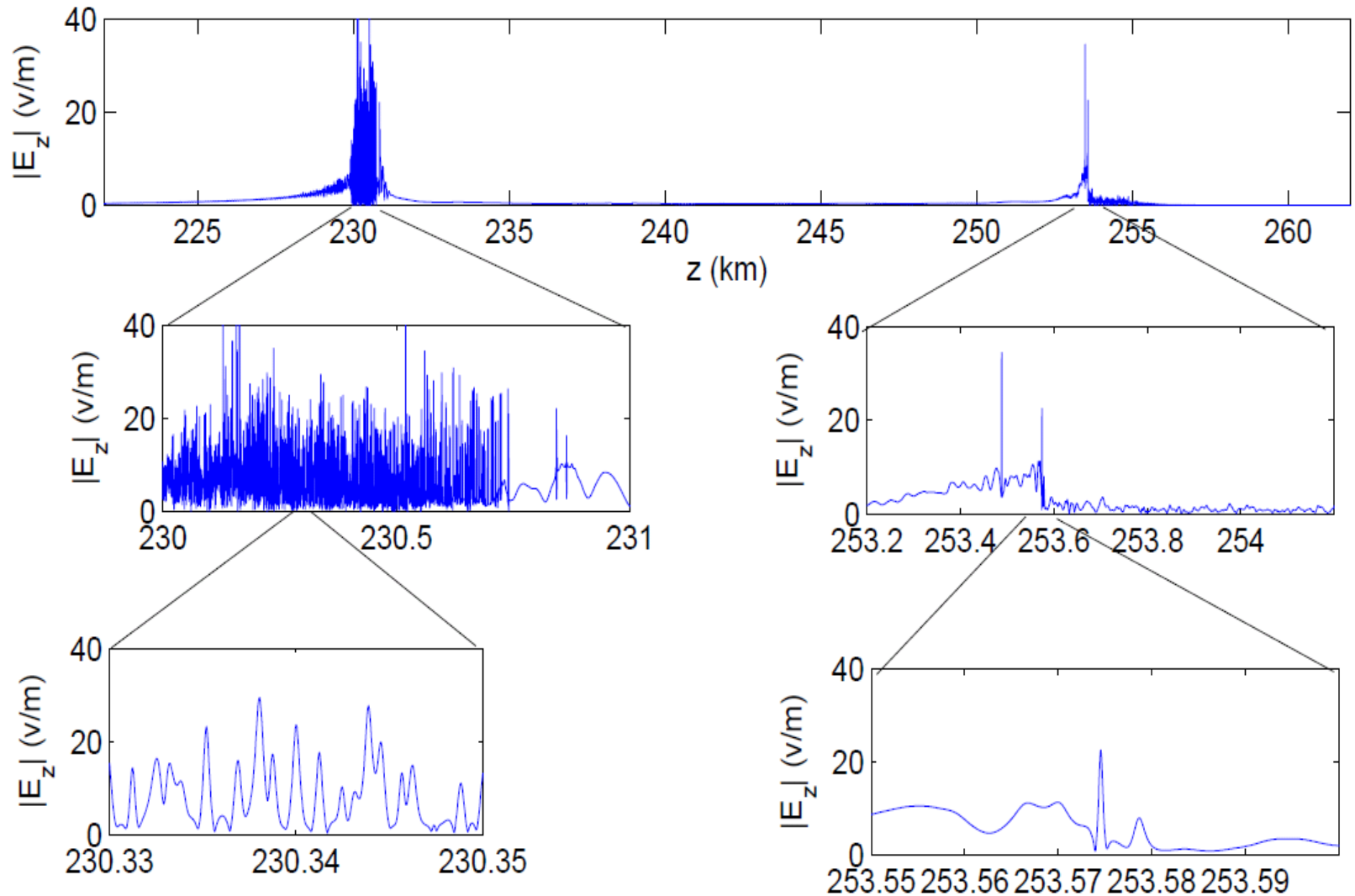
Simulation of topside ionospheric turbulence [Eliasson, 2008, 2010]

- A full scale model of the ionosphere is applied, in which O-mode EM wave, is injected from the bottom side. 1D ionospheric density profile has a Gaussian shape, B-field is tilted 14 deg. to vertical.
- The mathematical model includes equations for all three E-field components coupled with electron and ion continuity and momentum 1D equations.

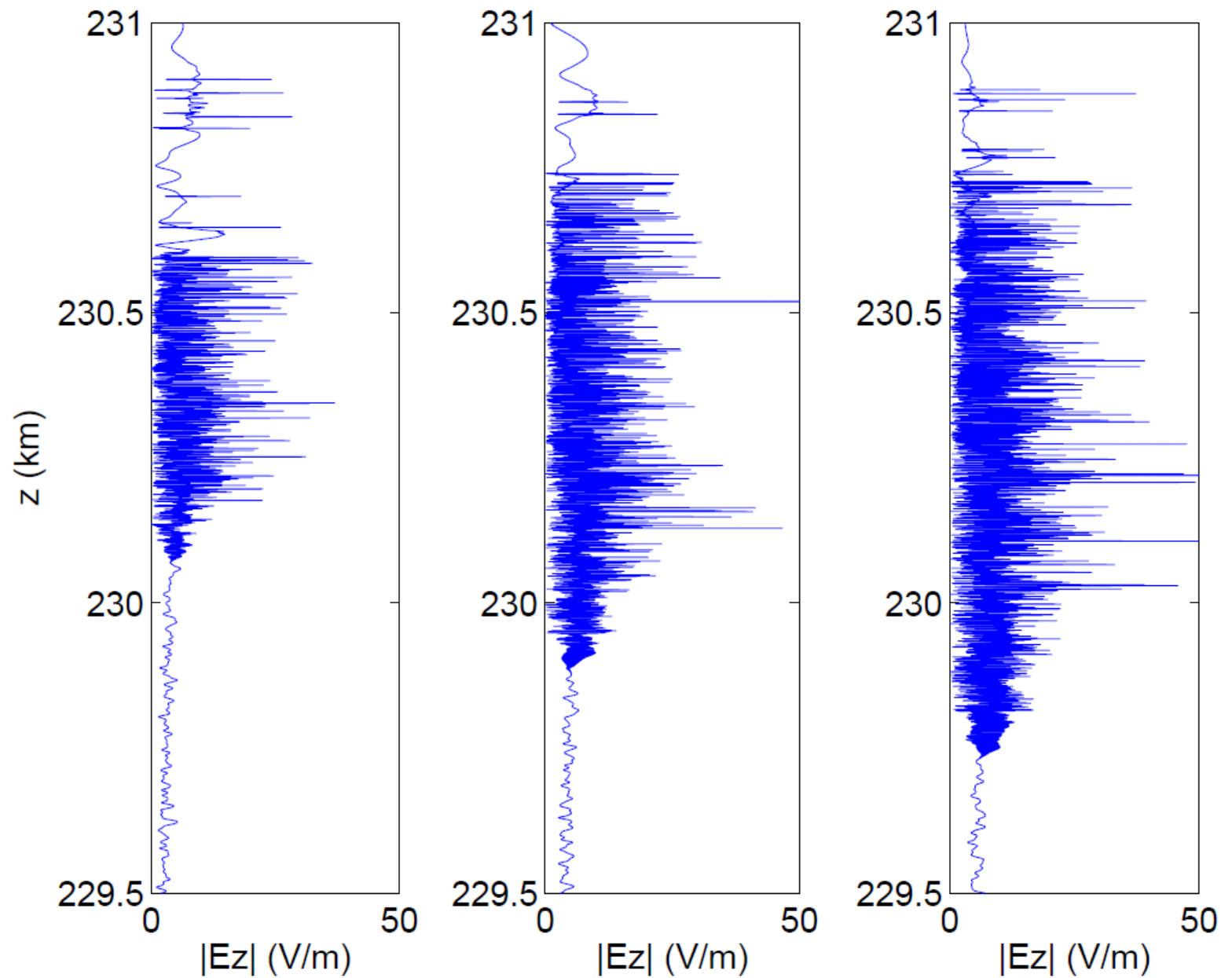




Most notably is the strong E_z field at ~ 230 km which is associated with electrostatic Langmuir waves that have collapsed into solitary wave packets.



Amplitude of E_z -field at different heights, for injected $E_0=1.5$ V/m, O-mode.



E_z at the vicinity of O-mode cutoff, Injected are the waves of 1.0 V/m, 1.5 V/m and 2.0 V/m.

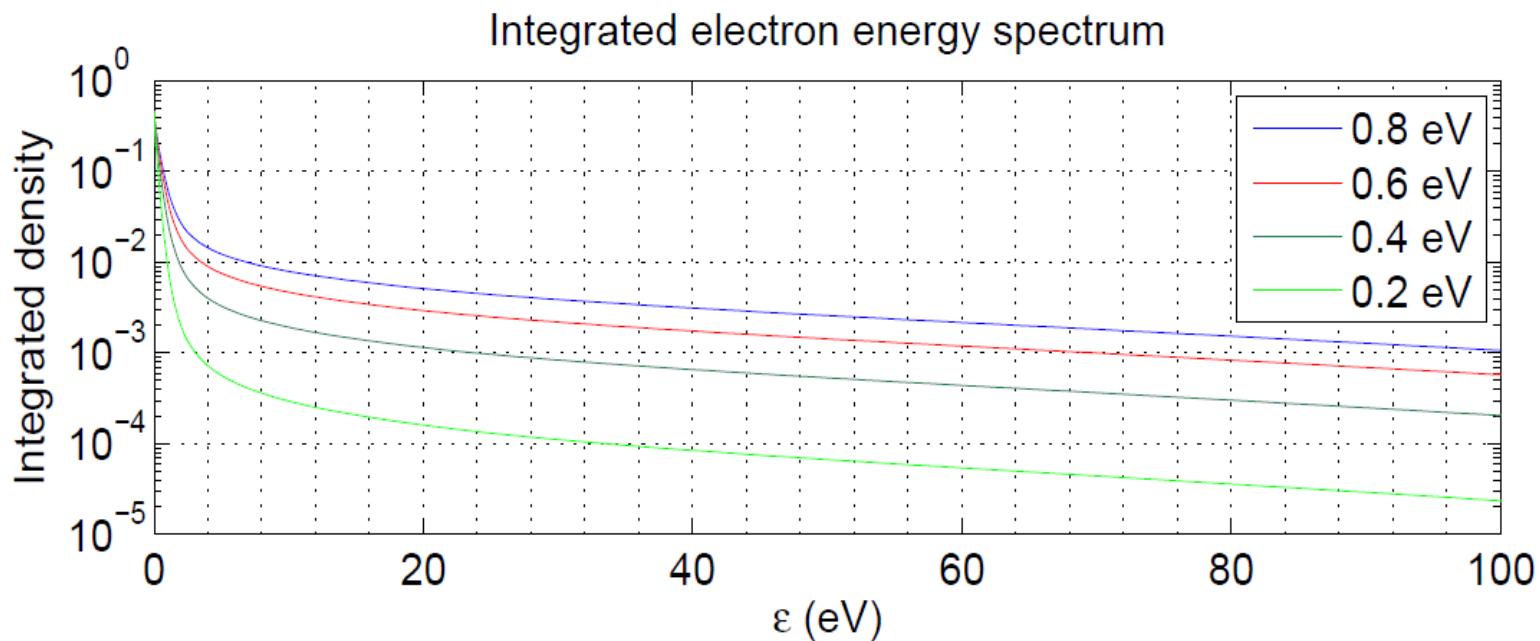
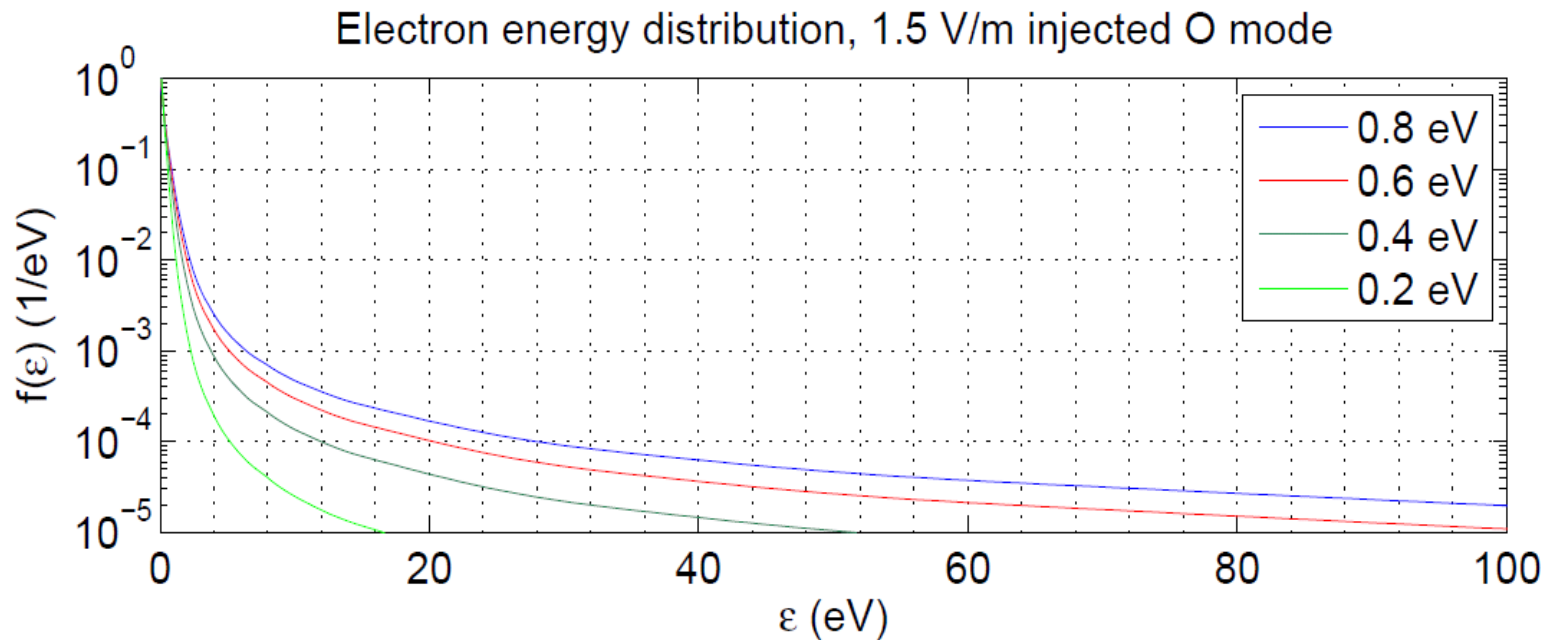
Electron acceleration by quasi-linear diffusion by plasma waves

Slowly varying electron distribution function is governed by F-P equation

$$\frac{\partial f}{\partial t} + v \frac{\partial f}{\partial z} = \frac{\partial}{\partial v} D(v) \frac{\partial f}{\partial v}$$

$$D(v) = \frac{\pi e^2}{m^2} \int W_k(\omega, k) \delta(kv - \omega) dk$$

W_k is the spectral density of the electric field per wave number



The energy distribution and integrated density.

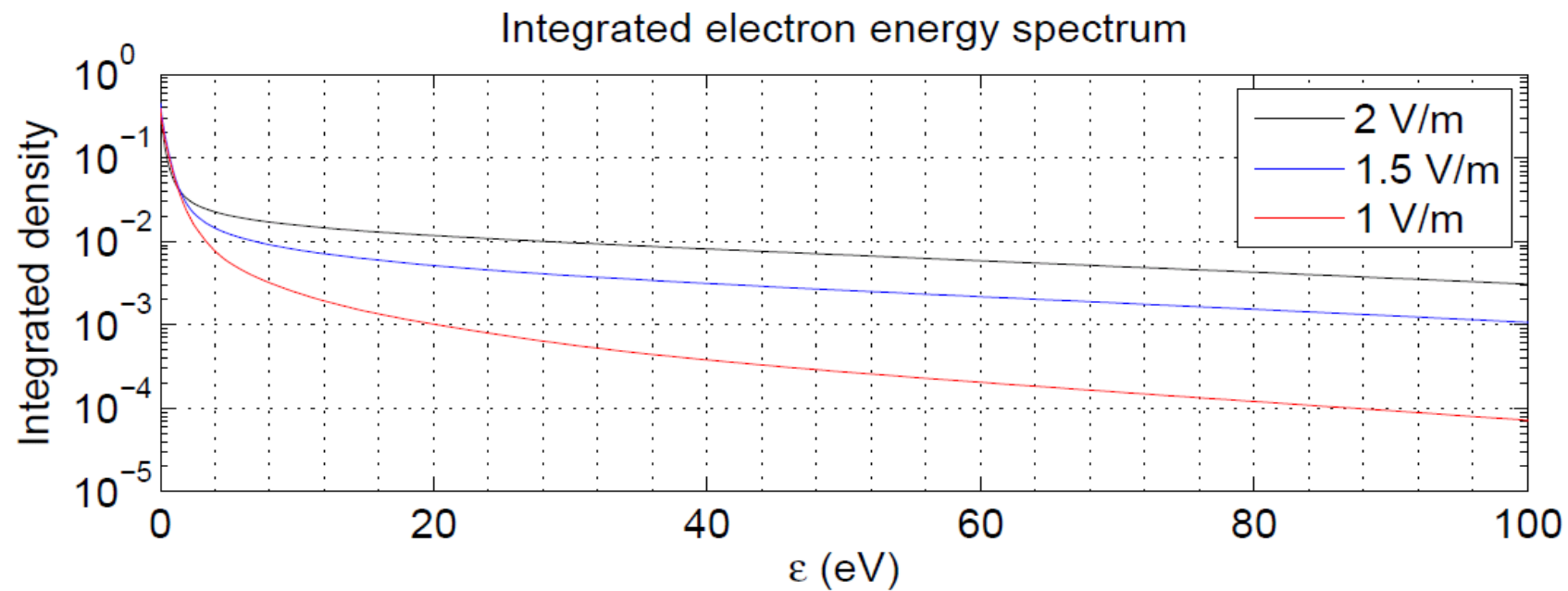
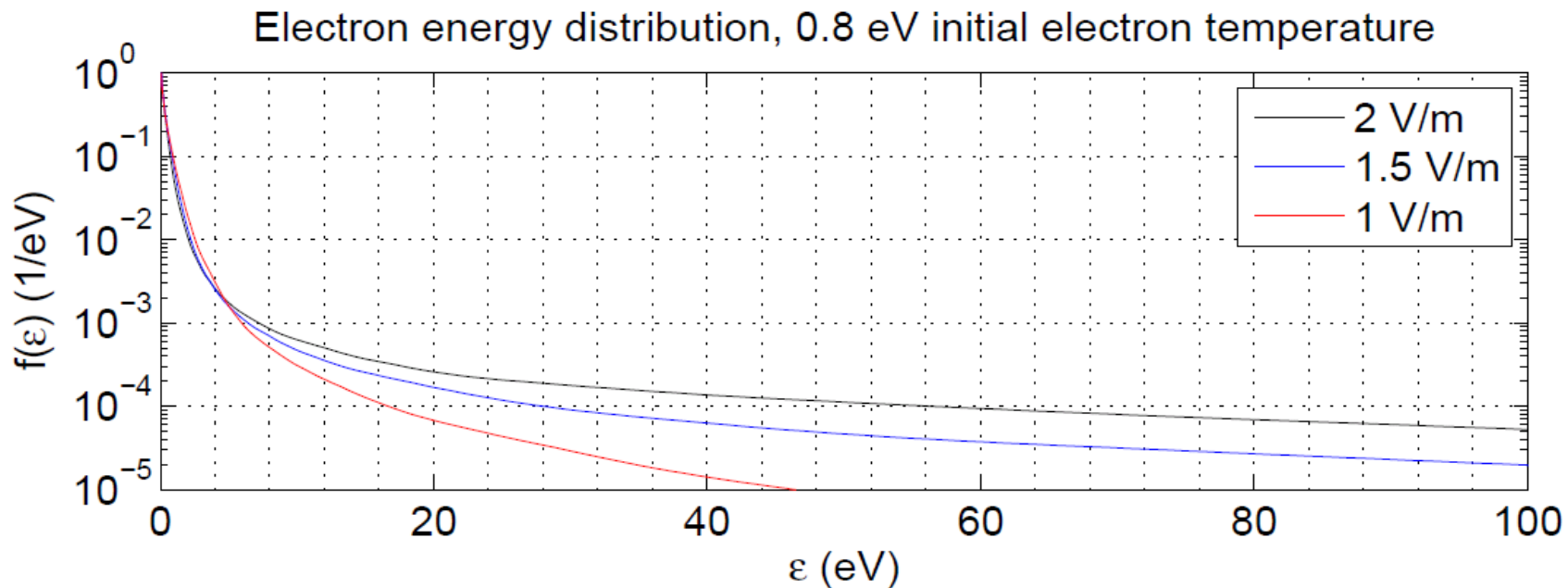
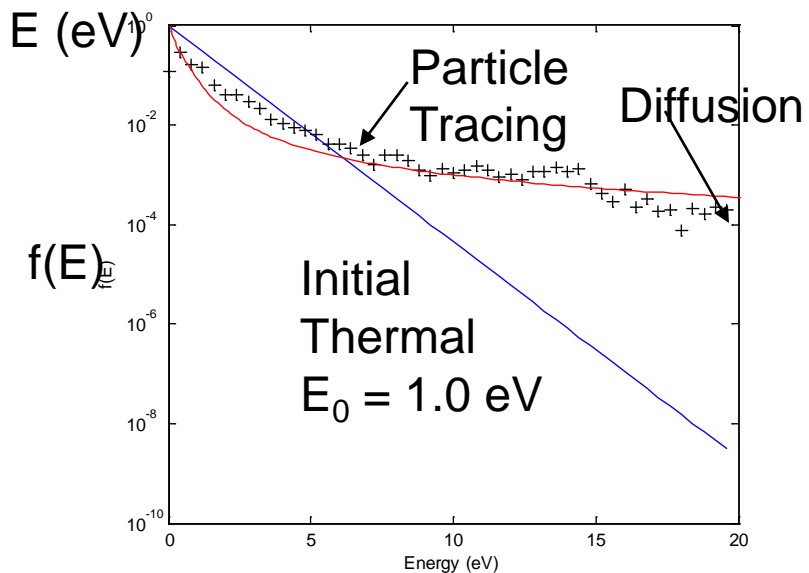
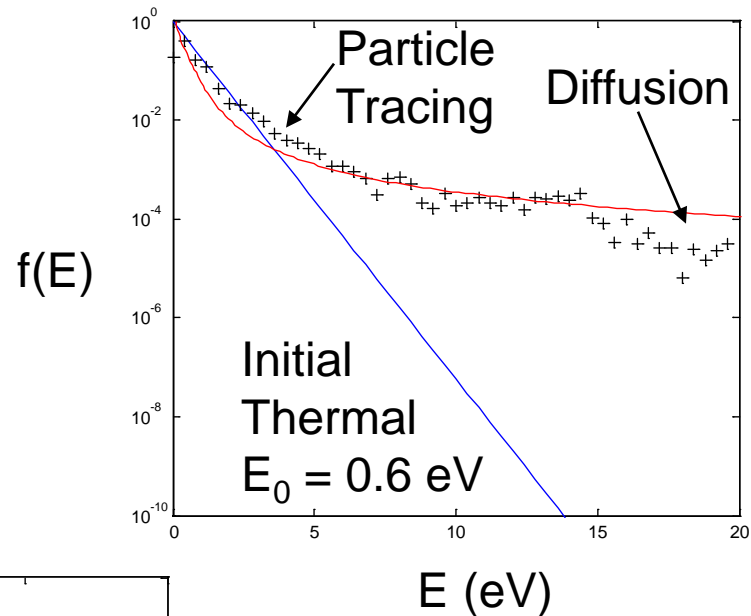
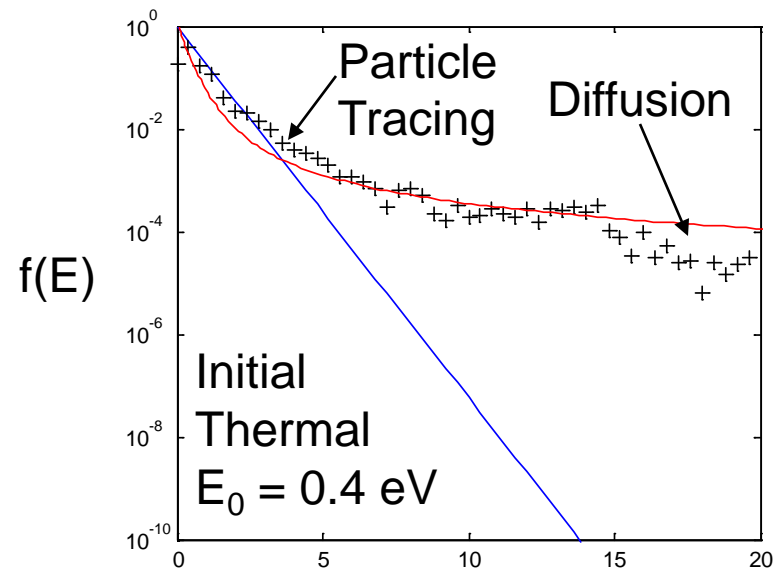


Fig. 20. The energy distribution and integrated density.

Comparing Diffusion Calculation with Particle Tracing



Modeling Flux Transport

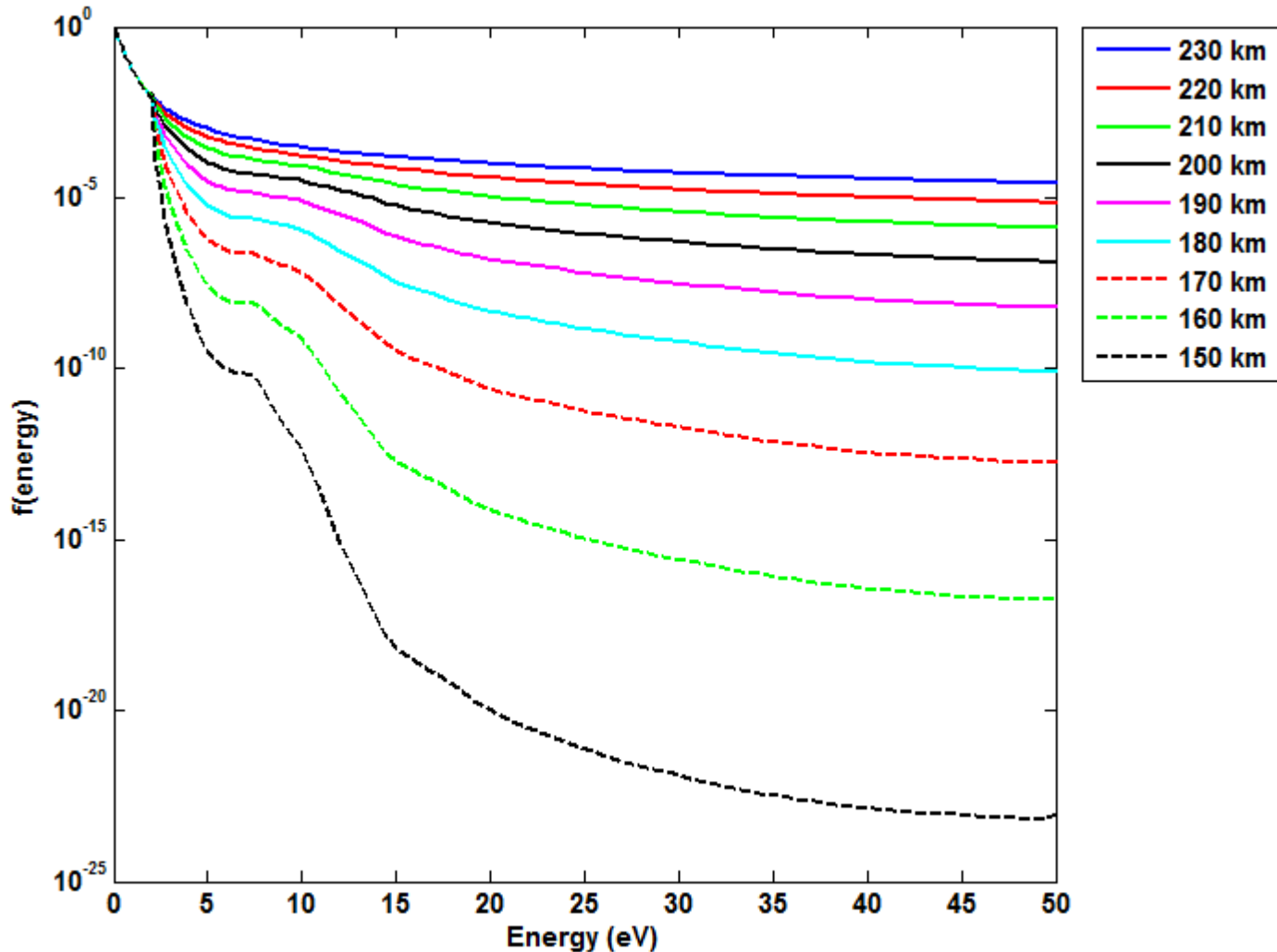
Gurevich et al. [1985] showed that propagation of fast electrons outside of accelerating layer is described by the kinetic equation:

$$\frac{d}{dz} \frac{v}{v_{tr}(z, \varepsilon)} \frac{df_0(z, \varepsilon)}{dz} - \frac{3v_{in}(z, \varepsilon)}{v} f_0(z, \varepsilon) = 0$$

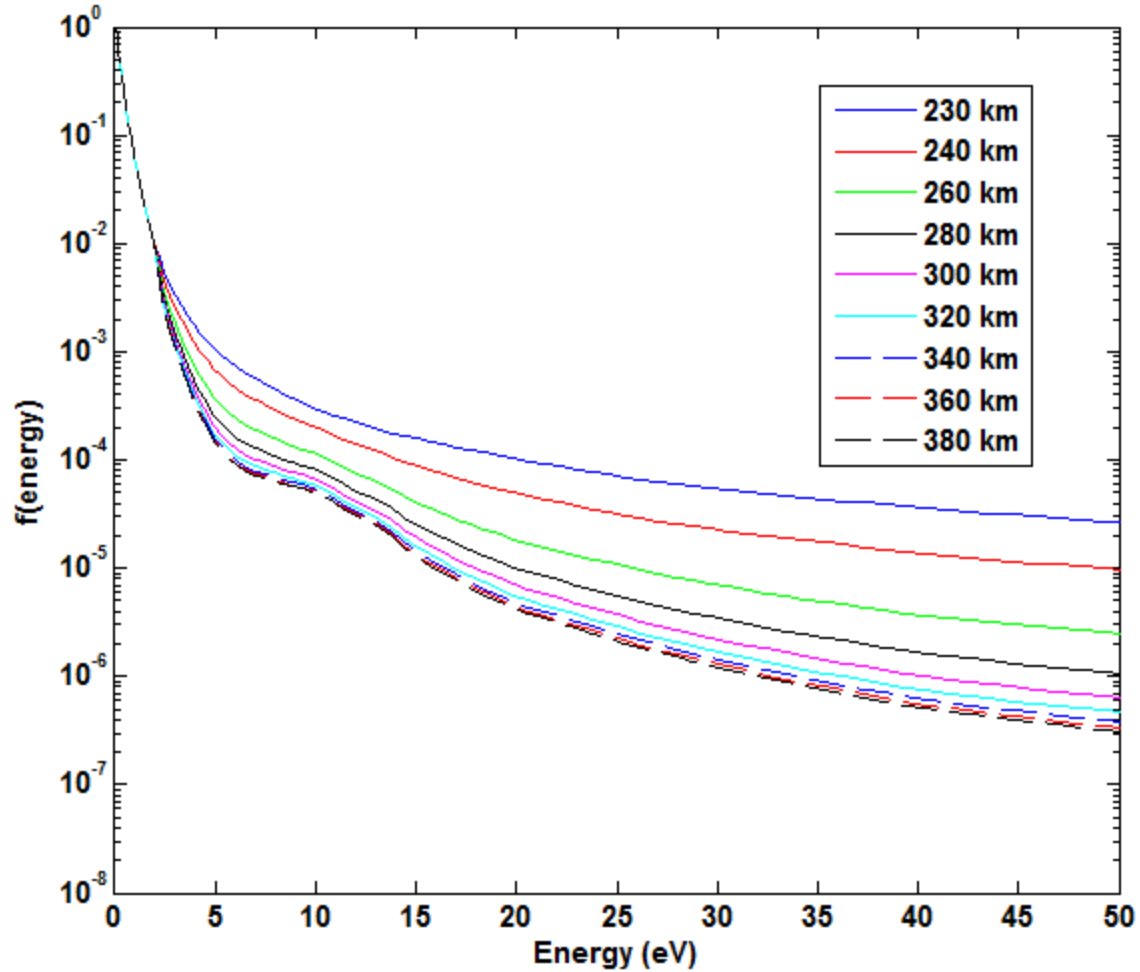
$$f(E, z) = f_0(E) \exp\left(-\int_{z_0}^z \frac{dz}{L(E, z)}\right)$$

$$L(E, z) = 1 / \left(\sqrt{3(N_O^2 \sigma_t^O \sigma_{in}^O + N_{N_2}^2 \sigma_t^{N_2} \sigma_{in}^{N_2} + N_{O_2}^2 \sigma_t^{O_2} \sigma_{in}^{O_2})} \right)$$

$E = 1.5 \text{ V/m}$, Simulation for HAARP, $E_0 = 0.6 \text{ eV}$,
 $Z_0 = 230 \text{ km}$, O Mode: $f(E)$ for $z = 150\text{-}230\text{ km}$



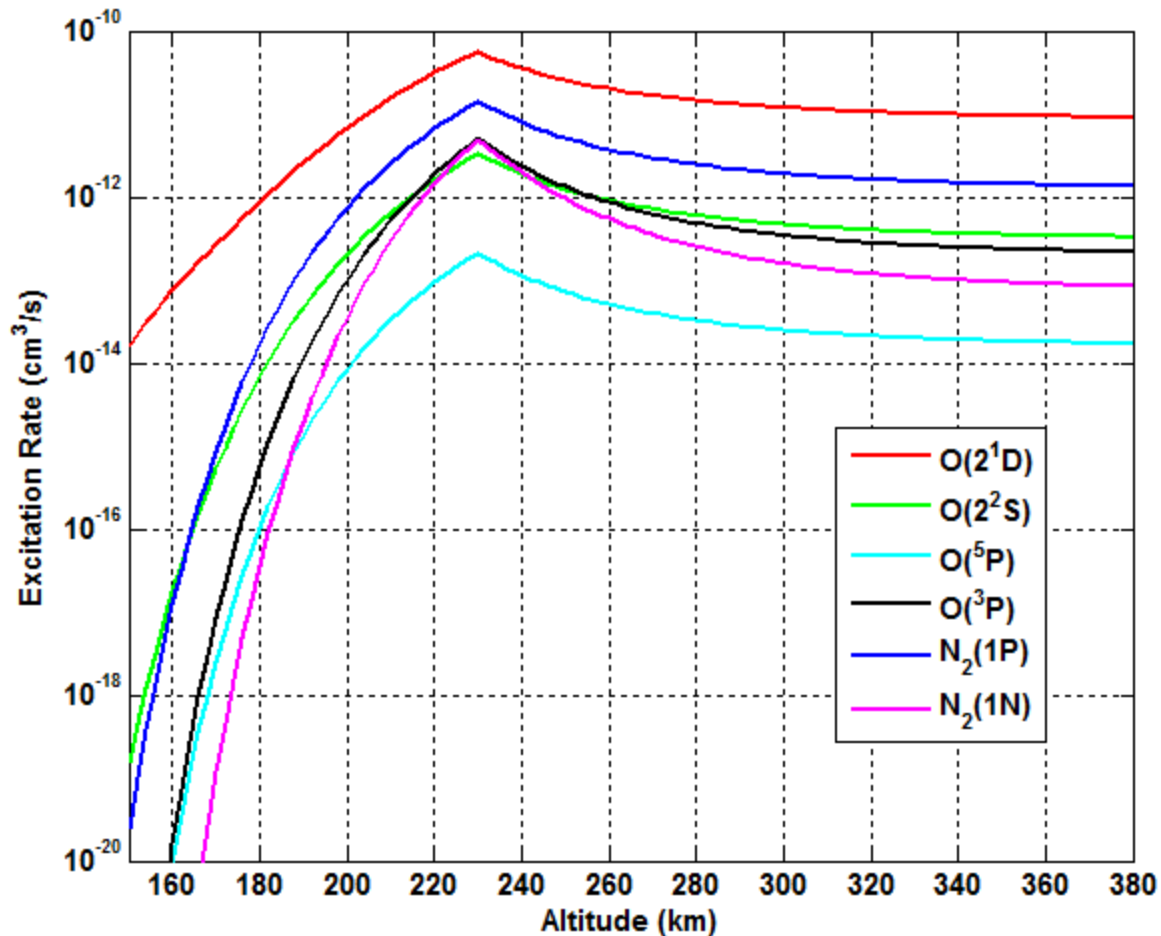
$E = 1.5 \text{ V/m}$, Simulation for HAARP, $E_0 = 0.6 \text{ eV}$,
 $Z_0 = 230 \text{ km}$, O Mode: $f(E)$ for $z = 230\text{-}380 \text{ km}$

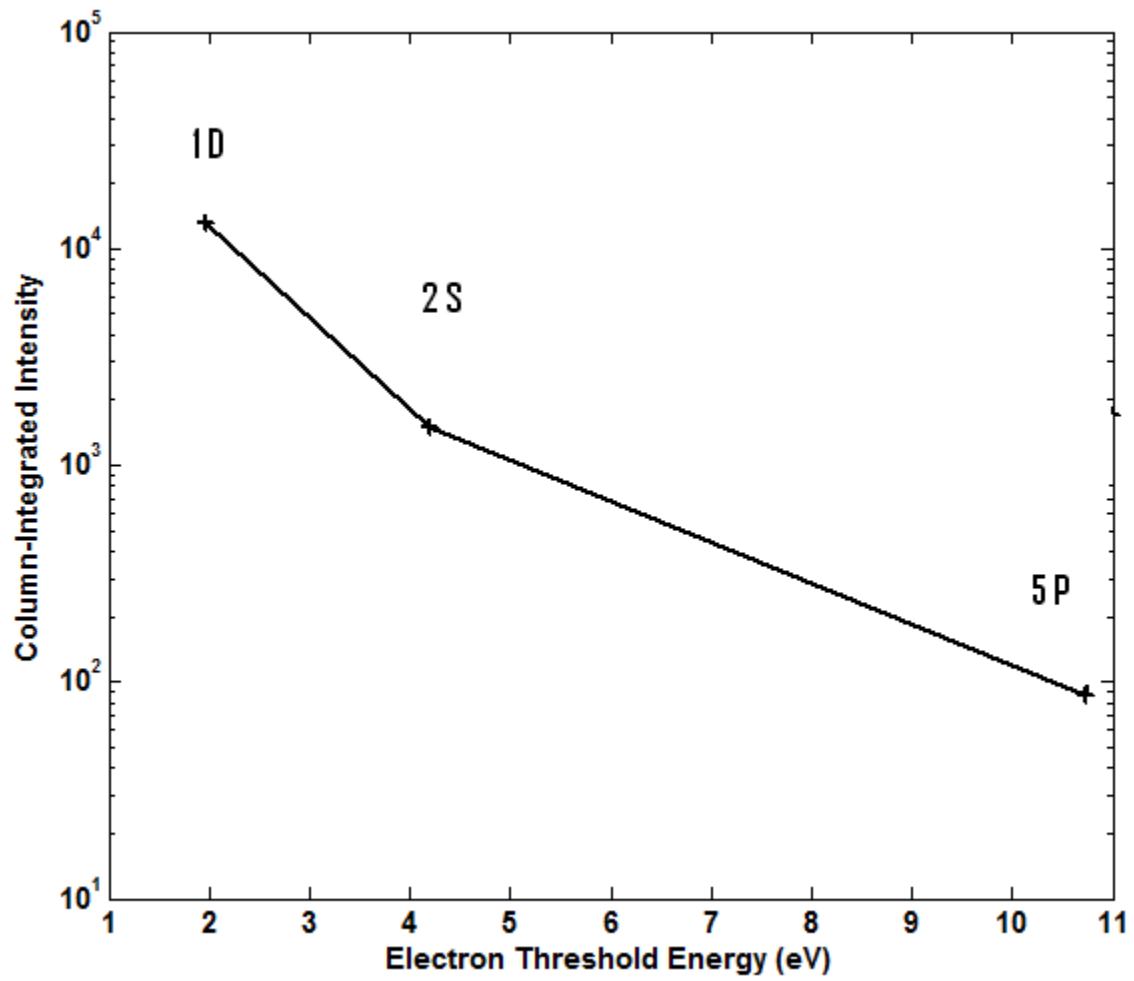


Optical emissions

Electronic Level	Emission, nm	Threshold, eV	Misc.
O(¹ D)	630	1.96	Red line
O(² S)	557.7	4.19	Green line
O(⁵ P)	777.4	10.74	
O(³ P)	844.6	10.99	
N ₂ (B ³ Π _g)	2,531-478	7.35	1P (645.4 nm band (8,5))
N ₂ (B ² Σ _u ⁺)	587-286	18.8	1N (391.4 nm band (0,0))

Modeling Optical Emission: Excitation Rate ($E = 1.5$ V/m, Simulation for HAARP, $E_0 = 0.6$ eV, $Z_0 = 230$ km, O Mode)

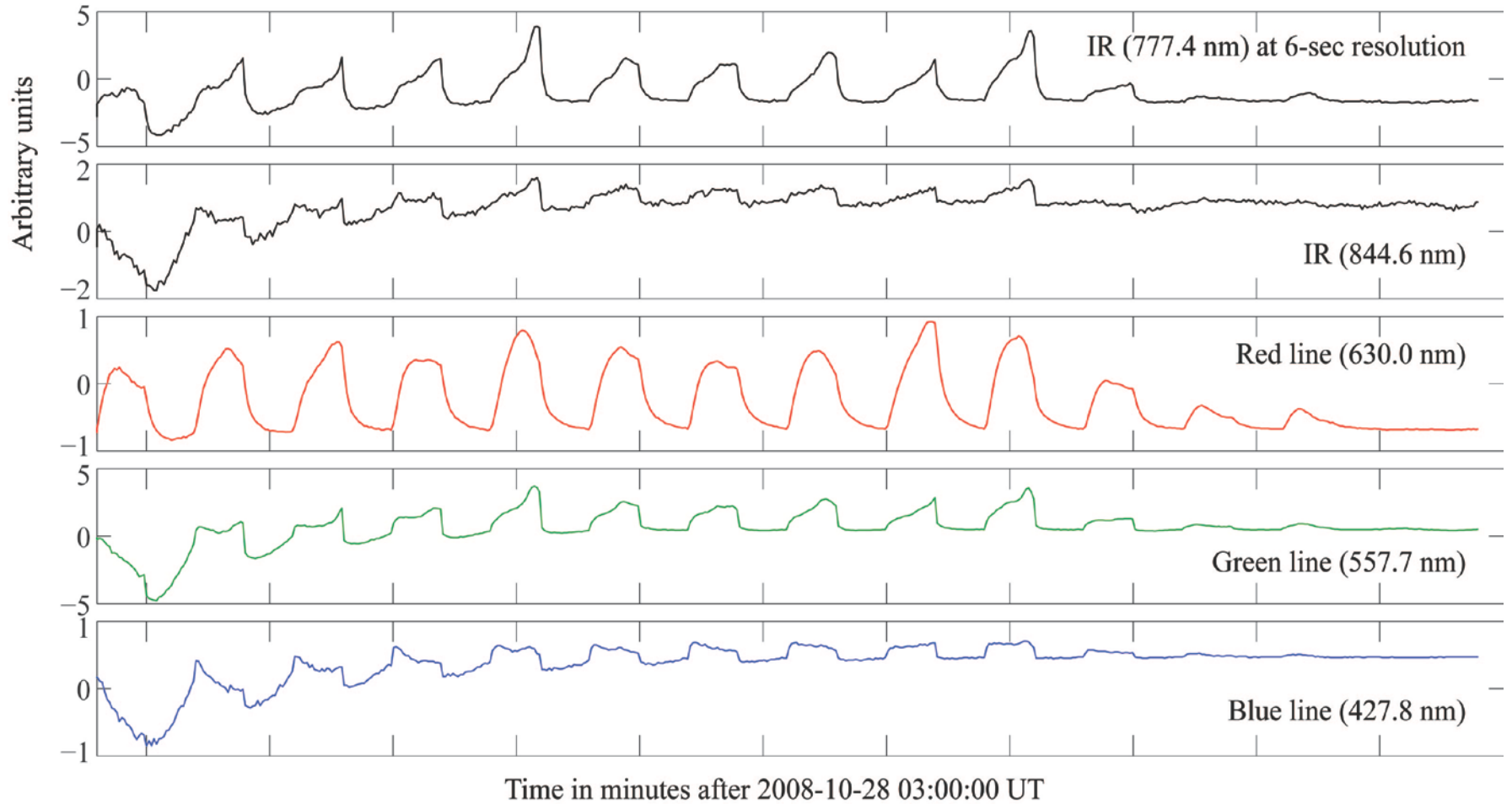




Model of the column integrated optical emissions of atomic oxygen

Telescopically-imaged heater-induced airglow

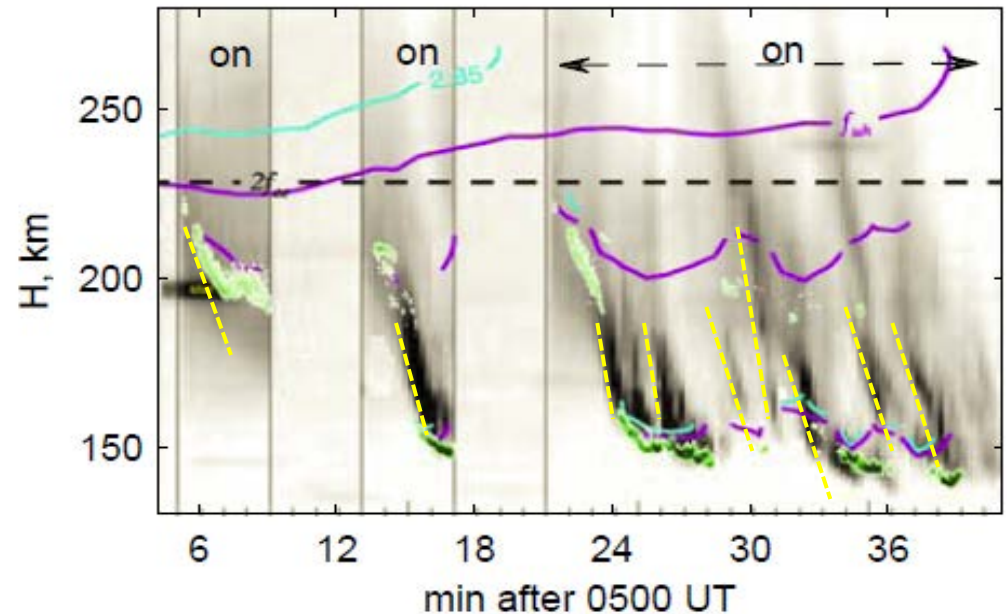
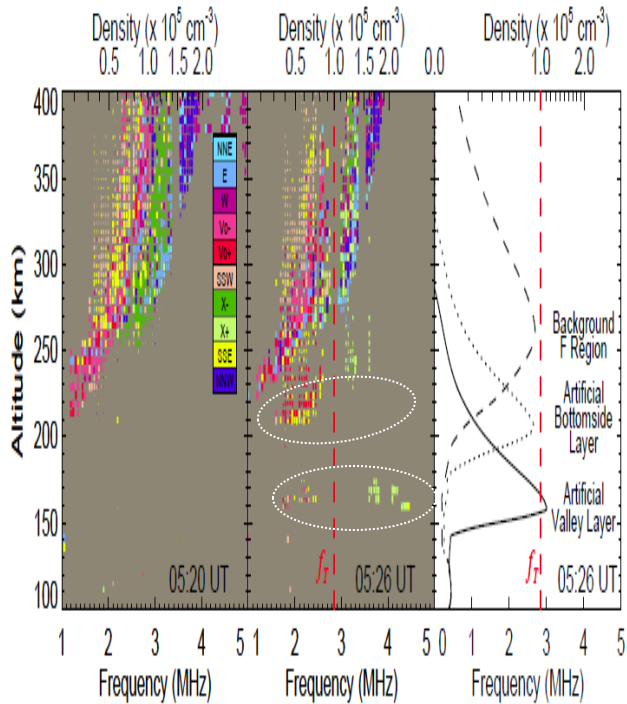
Stanford/HAARP Photometer



Formation of the ionized layers

Descending AA

2GH, 440 MW, MZ



- (left) Background echoes (the heater off).
- (center) Heater on: Two lower layers of echoes near 160 and 200 km virtual height for 210 s.
- (right) True height profiles.

Time-vs-altitude plot of **557.7 nm** optical emissions along **B** with contours showing the altitudes where $f_p = 2.85 \text{ MHz}$ (blue), $UHR = 2.85 \text{ MHz}$ (violet), and $2f_{ce} = 2.85 \text{ MHz}$ (dashed white). Horizontal blips are stars. **Shown in green is the Ion Acoustic Line intensity.**

✓ the artificial plasma near h_{min} was quenched several times.

Modeling the ionization wave

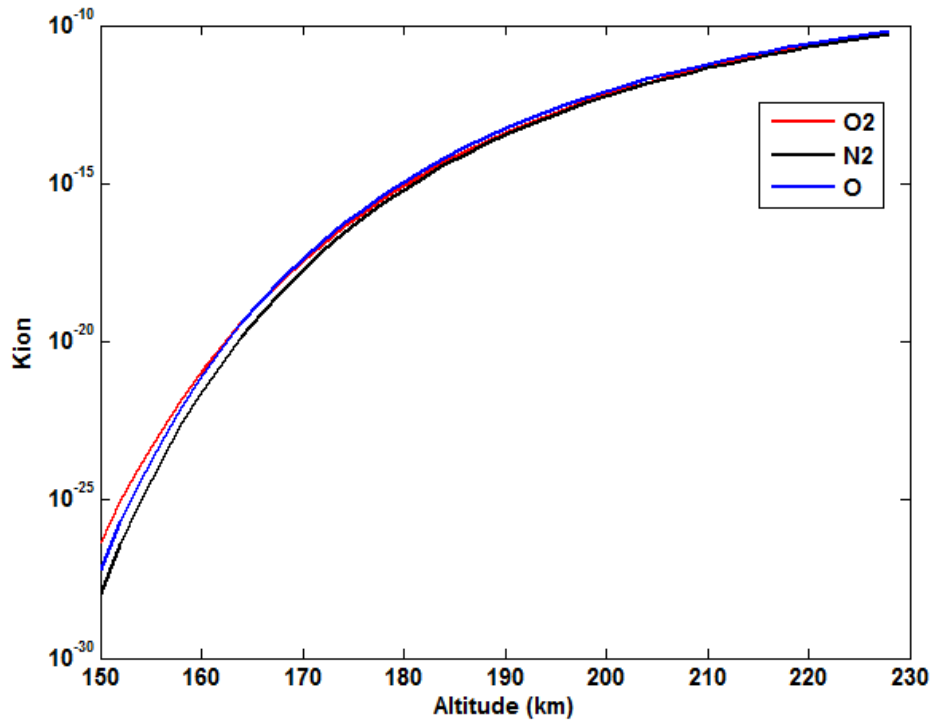
The ionization wave occurs in steps, the descending suprathermal electrons ionize the neutral species till the plasma reaches such density that it reflects the HF wave. The reflected wave generates the Langmuir turbulence, leading to the electron acceleration.

A 1D model of a uniform plasma slabs is applied:

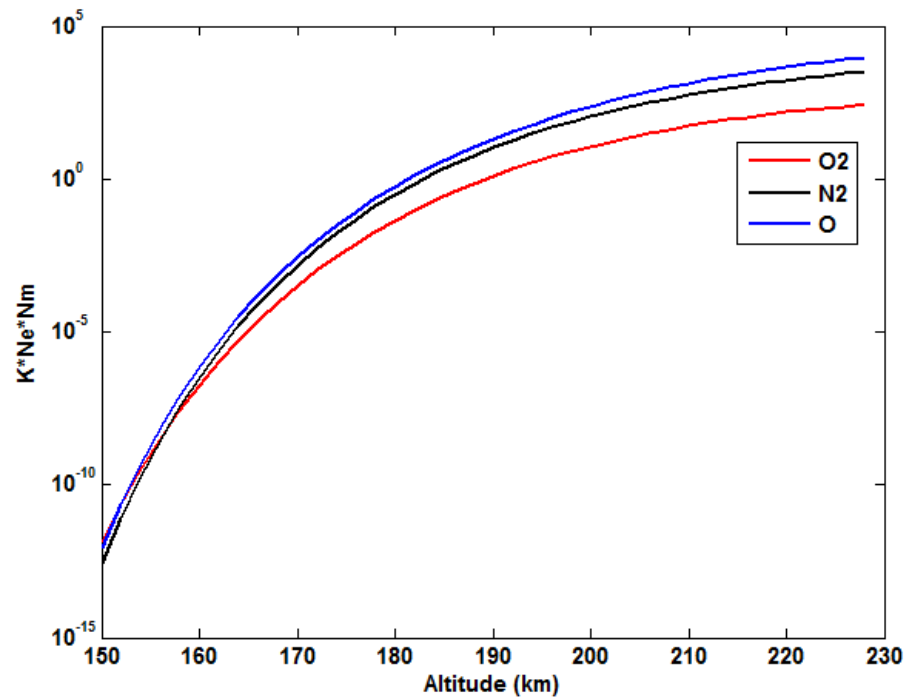
- The ionization is caused by descending flux of suprathermal electrons (energy losses included)
- E-I recombination, charge exchange and ambipolar diffusion all take place inside a slab

$E_0 = 0.6 \text{ eV}$, $E_{in} = 1.5 \text{ V/m}$

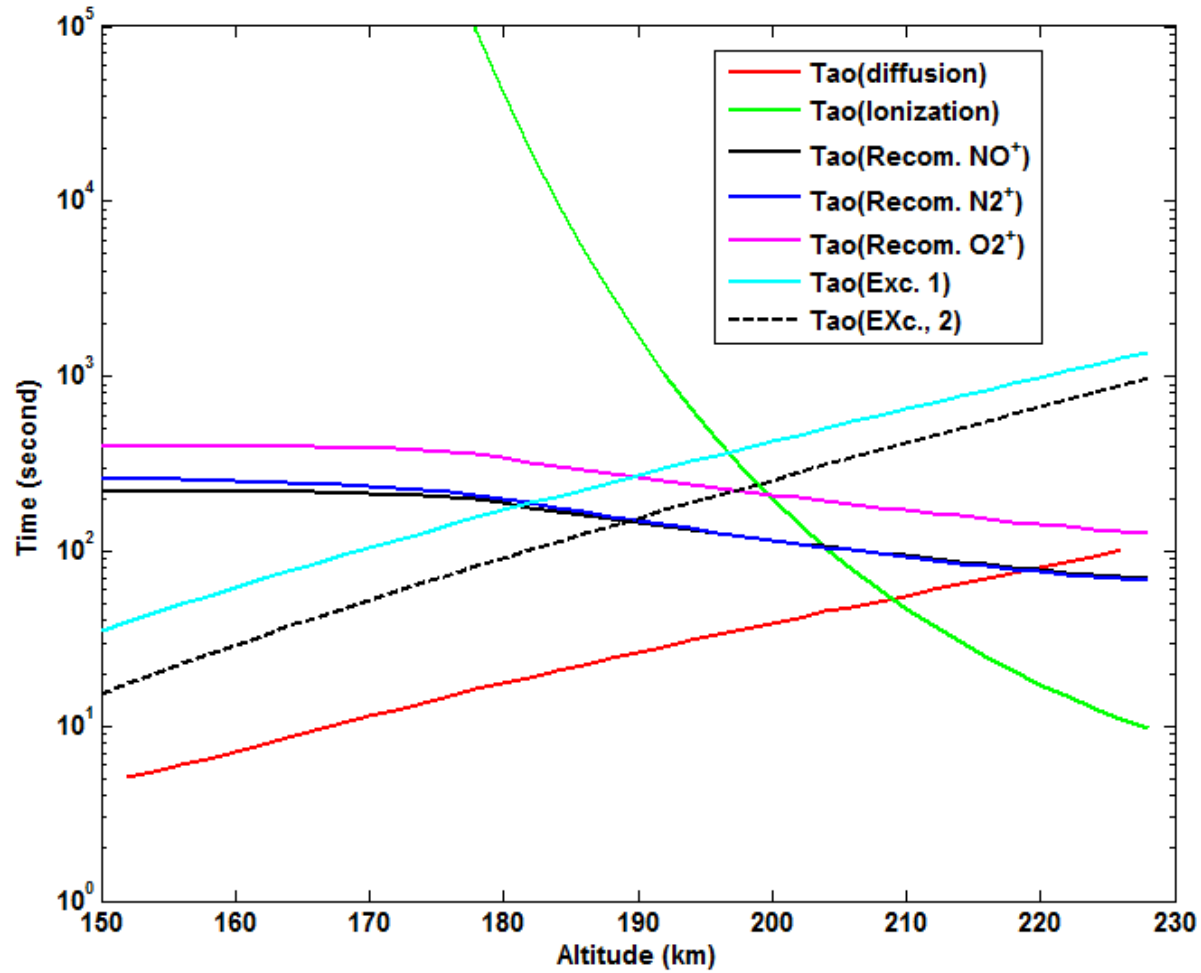
K_{ion}



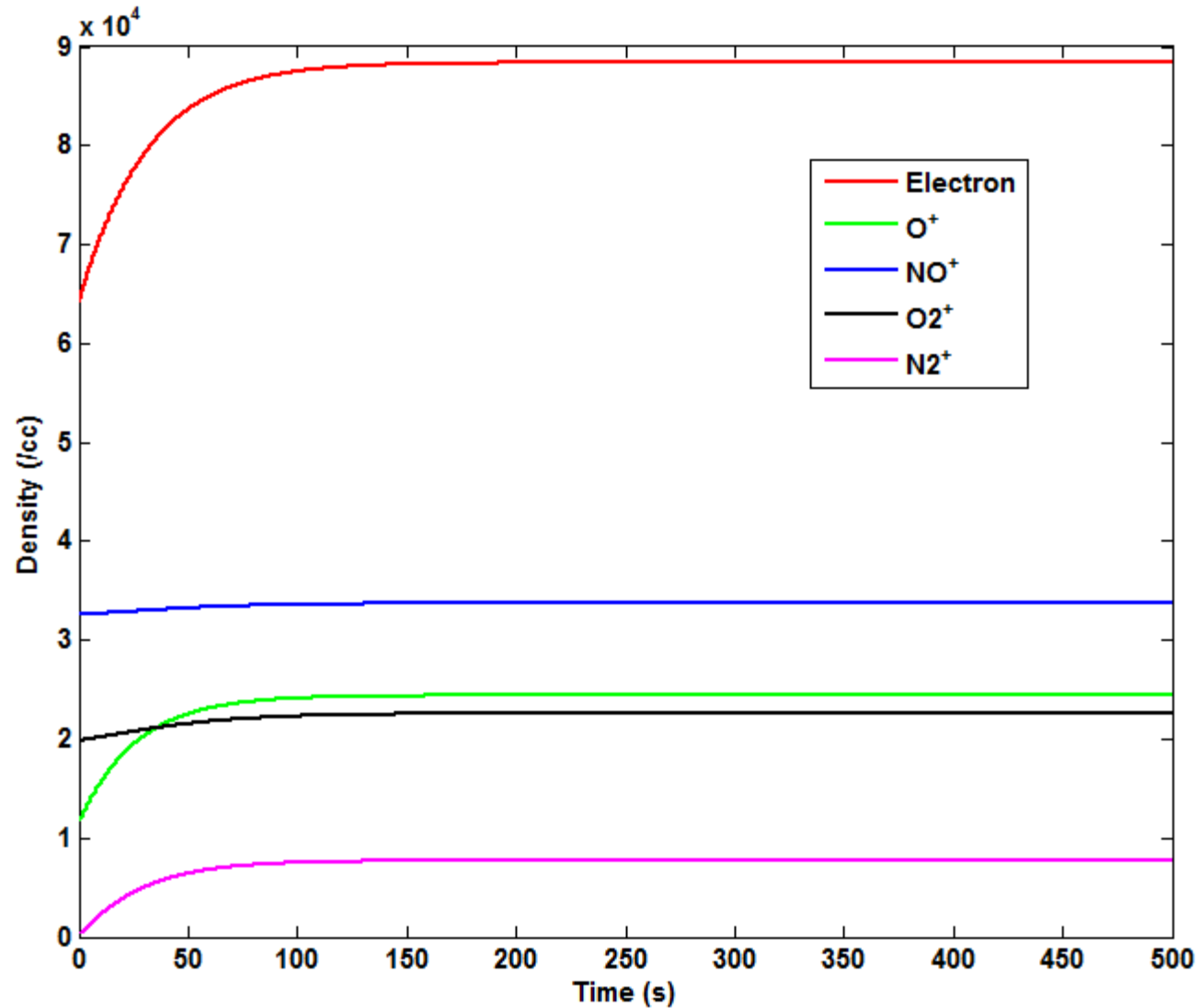
$K_{ion} * N_m * N_{e0}$



Timescales at 216 km

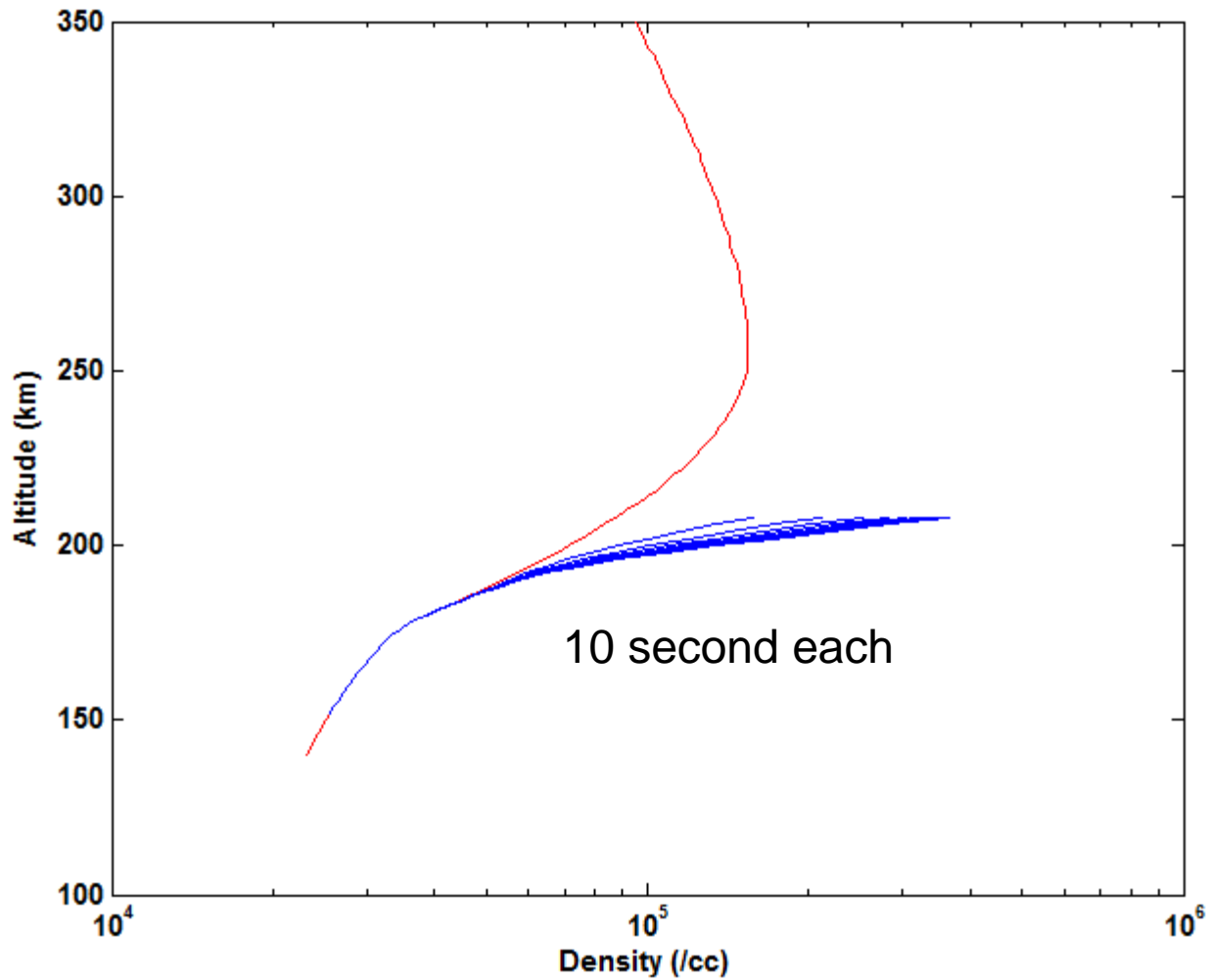


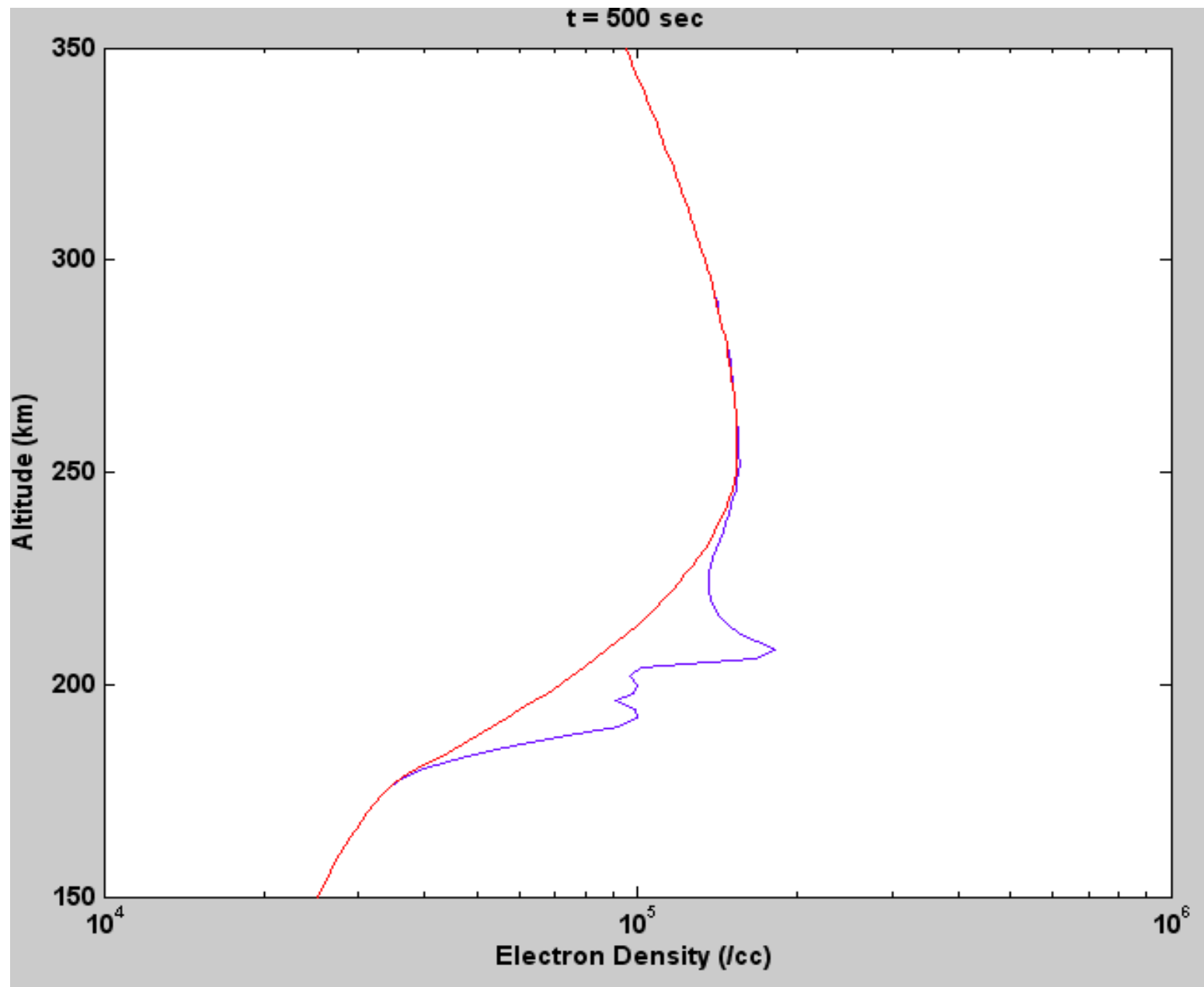
Electron and ion density evolution at 200 km



Electron Density Profile

Heating layer moved to $z = 216$ km





Conclusions

- A model of artificial optical emissions and formation of the ionized layers is discussed
- The model is based on the electron acceleration by the Langmuir turbulence caused by HF-heating
- The ionization wave which forms the ionized layers occurs in steps
- The model will be checked against the existing observations